

## Colloidal Probe Reflection Interference Contrast Microscopy (RICM)

### Description

The RICM technique was first introduced to study cell adhesion, in particular for the visualization of the area of the cell attached to the surface in the focal adhesion site Curtis, 1964. In the 1980s, Sackmann et al. demonstrated that the distances between a planar transparent substrate and spherical objects, like a colloidal bead, hovering over the substrate can be well quantified by RICM Kuhner & Sackmann, 1996 Radler & Sackmann, 1997 Schiling et al., 2004. We applied colloidal probe RICM to determine the thickness of the HA films modified by proteins. Moreover, tracking the normal and lateral displacement of HA functionalized beads hovering above a planar HA film provided insight into interactions, such as cross-linking in and between HA assemblies.

**The working principle** RICM is based on the interference of light reflected from different interfaces (Fig. 3).

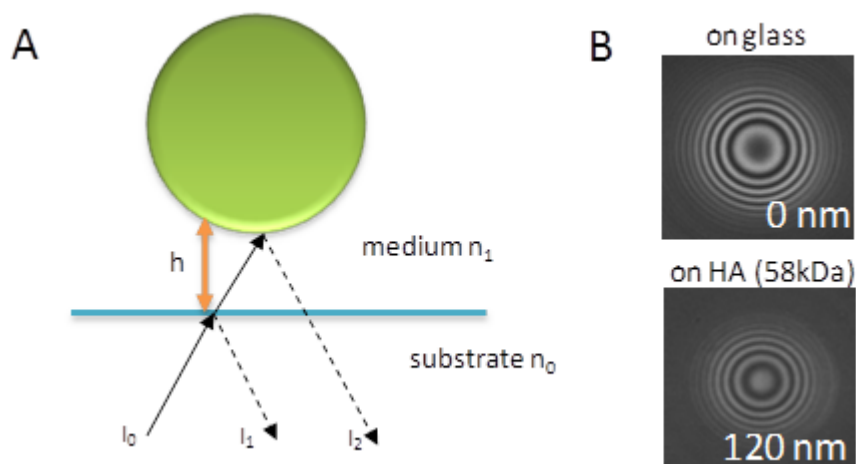


Figure 3. Schematic representation of the RICM working principle. A. The incidence beam with intensity  $I_0$  is partly reflected from the substrate interface (bare glass or glass covered with a thin transparent gold layer) and partly from the surface of a colloidal bead. The reflected beams with intensities  $I_1$  and  $I_2$  interfere and give rise to an interference pattern. B. Interferographs of a colloidal bead sitting on the glass surface or on an HA-film of

approximately 120 nm thickness, respectively.

The incidence light is first partly reflected at the glass/medium interface. The transmitted light is further partly reflected at the colloidal probe/medium interface. The reflected beams interfere and give rise to a pattern, so-called Newtonian rings of maximum ( $I_{max}$ ) and minimum intensity ( $I_{min}$ ). The relation between the intensity distribution and height of the bead above the surface is given by:

$$\text{Eq. 8} \quad I(h(x,y)\lambda) = I_{min} + I_{max} + 2\sqrt{I_{min}I_{max}}\cos\left[\frac{4\pi n_d(x,y)}{\lambda} + \delta\right]$$

where  $h(x,y)$  is the distance at a given lateral position with coordinates  $x$  and  $y$ ,  $n_b$  is the refractive index of the medium,  $\delta$  is the phase shift of the reflected light and  $\lambda$  is the wavelength.

**Reconstruction of the film thickness** For a colloidal probe of spherical geometry, the interference pattern consists of circular fringes. The intensity in the center changes as the distance between probe and substrate varies (Fig. 4A).

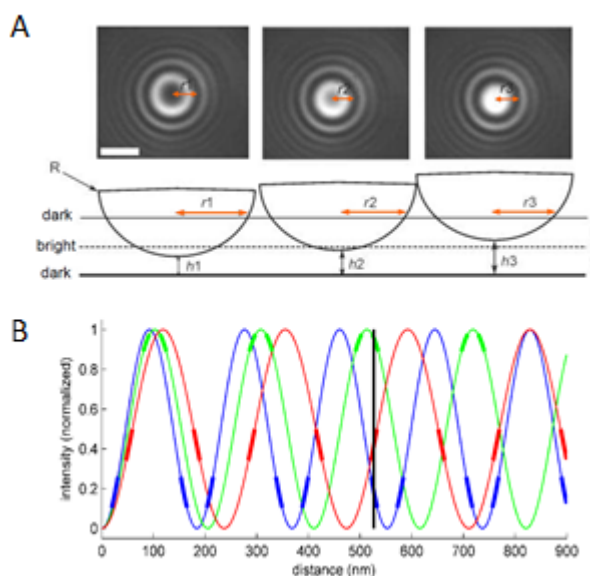


Figure 4. Quantification of the film thickness from RICM images. A. Height reconstruction. RICM fringes showing the fringe pattern of a colloidal bead of different heights from the surface ( $h_1$ ,  $h_2$  and  $h_3$ ). The intensity in the center changes as a function of the height, and the radius of first fringe changes accordingly. B. Normalized intensities in the center of the probe as a function of probe-substrate distance, for three selected wavelengths 490nm, 546 nm and 630 nm (triple-wavelength RICM, thin colored lines), A given intensity, measured for a given wavelength (indicated by thick colored lines), would be consistent with various probe-substrate distances ("candidate distances") at each individual wavelength. By comparing different wavelengths, many candidate distances can be discarded. A unique distance remains that is consistently found with all three wavelengths. This approach is used to reconstruct distances above 100 nm. Copyright: Figures reproduced from Limozin, L. & Sengupta, K. *ChemPhysChem* 10, 2752-2768 (2009) and from Richter, R.P., et al., *Journal of the American Chemical Society* 129, 5306-+ (2007)

The radius  $r_l$  from the center at which an extremum (maximum or minimum) occurs can be calculated with good accuracy from the radially averaged intensity profile. The probe-substrate distance can then be derived from the radial position of the selected extremum:

$$\text{Eq. 9} \quad h(r_l) = \frac{\lambda l}{2n} - R + \sqrt{R^2 - r_l^2} \quad \text{for maximum}$$

and

$$\text{Eq. 10} \quad h(r_l) = \frac{\lambda(2l + 1)}{4n} - R + \sqrt{R^2 - r_l^2} \quad \text{for minimum}$$

where  $R$  is the radius of the bead, and  $l$  is the fringe order (which needs to be appropriately chosen).

**Dual-/Triple –wavelength RICM** From equations 9 and 10, it is obvious that the distance measured by RICM is not unambiguously defined. The RICM intensity response as a function of probe-substrate distance is periodic (Fig. 4B), and several solutions are possible. Under the experimentally relevant conditions, the periodicity is on the order of a few 100 nm, but depends on the wavelength. Correlation of RICM data at multiple wavelength, therefore, enable unambiguous distance determination over a distance range that is much larger than the periodicity (Fig. 4B). With triple-wavelength RICM the absolute height can be determined over a distance range of more than 1 micrometer with an accuracy of a few nanometers Attili,et al., 2012Richter et al., 2007Schiling et al., 2004Attili & Richter, 2012.

**Mean squared displacement** The thermal motion of the colloidal bead hovering above the HA film in the  $x$ ,  $y$  and  $z$  directions can be quantified by RICM with nm resolution. The thermally driven motion of the colloidal probe hence becomes detectable by microscopy. When unconstrained, the probe performs a random walk in  $x$  and  $y$  due to Brownian motion. The average of the squared distance which the bead has moved in, for example, the  $x$  direction within a given time step  $\tau$ , the so-called mean square displacement (MSD), is given by:

$$\text{Eq. 11} \quad MSD = \langle x^2(\tau) \rangle = \langle |x(t) - x(t - \tau)|^2 \rangle$$

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Figure 5. Colloidal bead undergoing thermal motion. The track shows a typical random walk of HA-functionalized polystyrene beads of about 25  $\mu\text{m}$  diameter hovering above a HA

film. A. Movement in x and y directions. B. Movement in z direction. The length of the track corresponds to 100 data points recorded at 8 frames per second.

The MSD is a measure of the mobility of the probe, and can be correlated to diffusion. (Fig. 5). Sengupta et al. used the calculation of the MSD (in z direction) by colloidal probe RICM to characterize the viscoelastic properties of HA films Sengupta et al., 2003

Schematic representation of the RICM working principle. A. The incidence beam with intensity  $I_0$  is partly reflected from the substrate interface (bare glass or glass covered with a thin transparent gold layer) and partly from the surface of a colloidal bead. The reflected beams with intensities  $I_1$  and  $I_2$  interfere and give rise to an interference pattern. B. Interferographs of a colloidal bead sitting on the glass surface or on an HA-film of approximately 120 nm thickness, respectively.

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